

# 321D: the case of entrainment

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# Keele is Not Kiel (Germany) But Where is it?

West Midlands:



Keele area

is famous for pottery: Wedgwood, ...

Exciting [HyDeploy.co.uk](http://HyDeploy.co.uk) / SEND projects



# Plan

- 321D concept

- Entrainment:

  - 3D results

  - Theory

  - 1D implementation

  - 1D results: MS width, core masses, ...

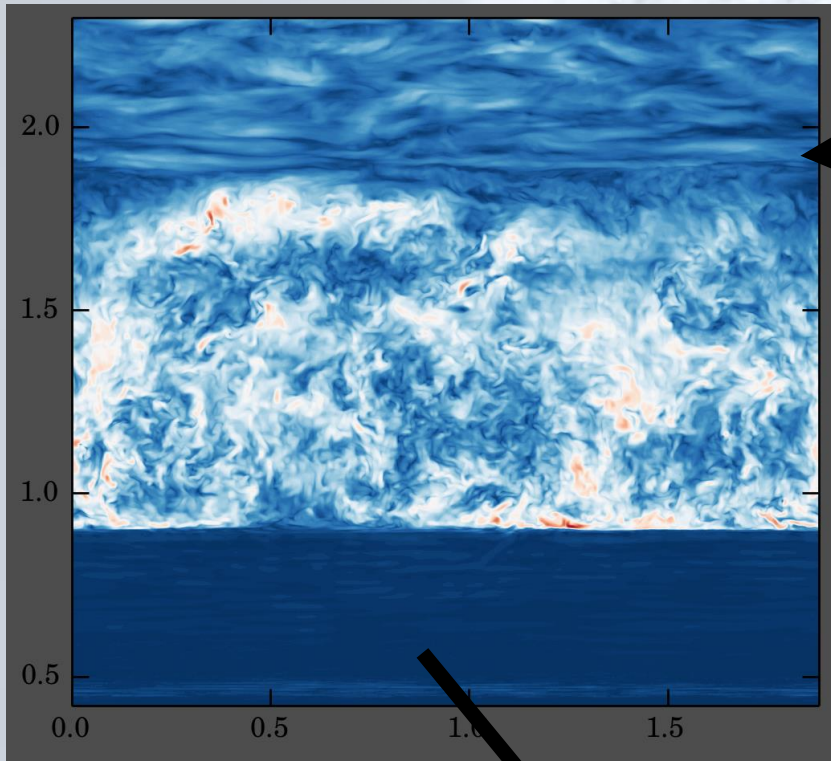
- Boundary shape - shear

- Conclusions & Outlook

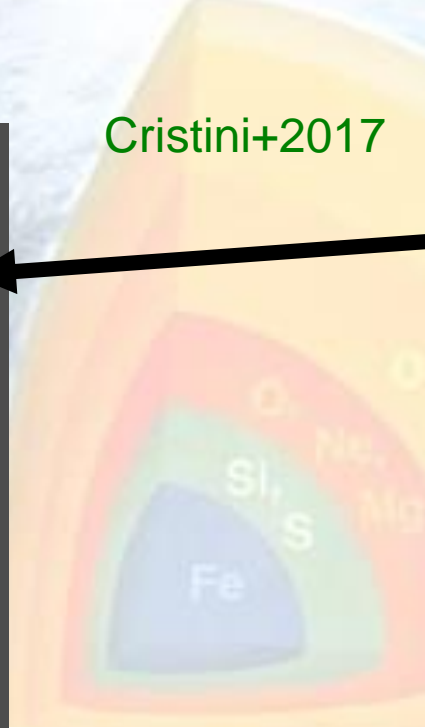


# Way Forward: 1 to 3 to 1D link

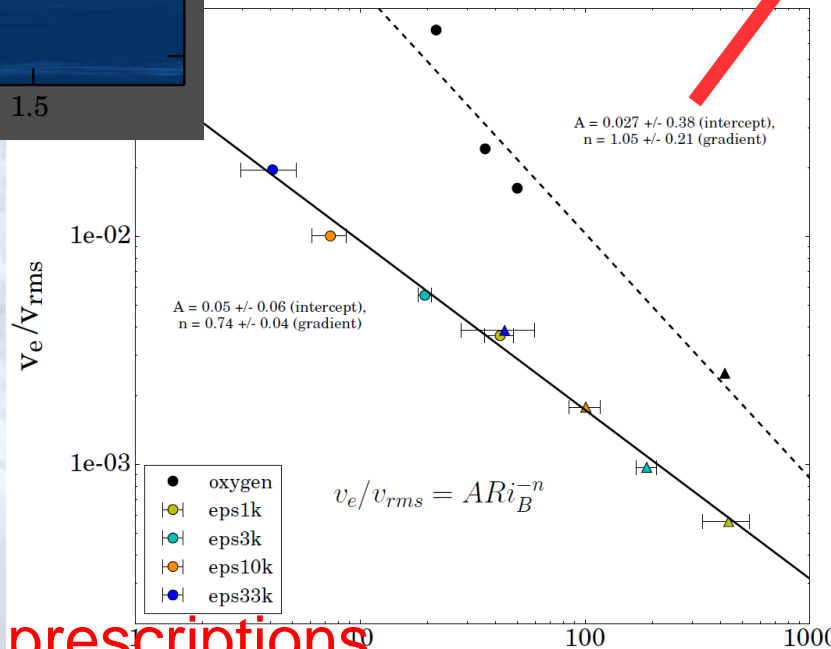
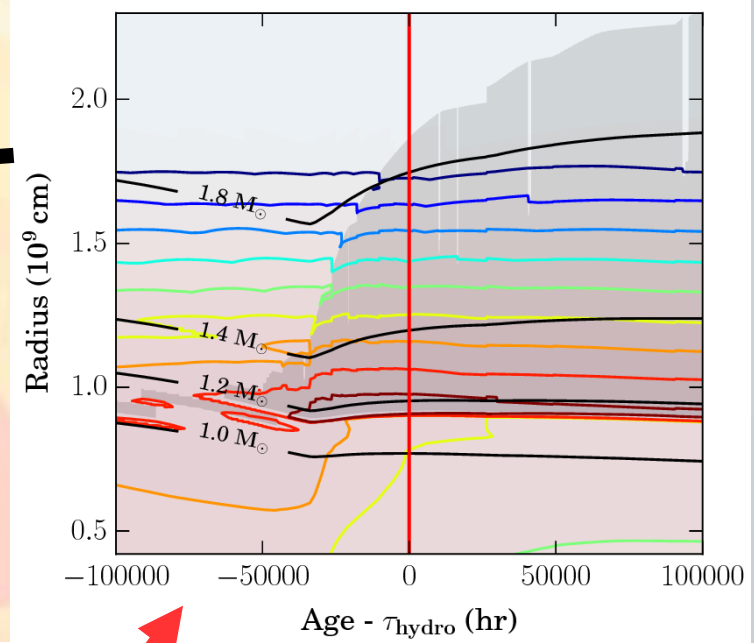
Targeted 3D simulations



Cristini+2017



Uncertainties in 1D

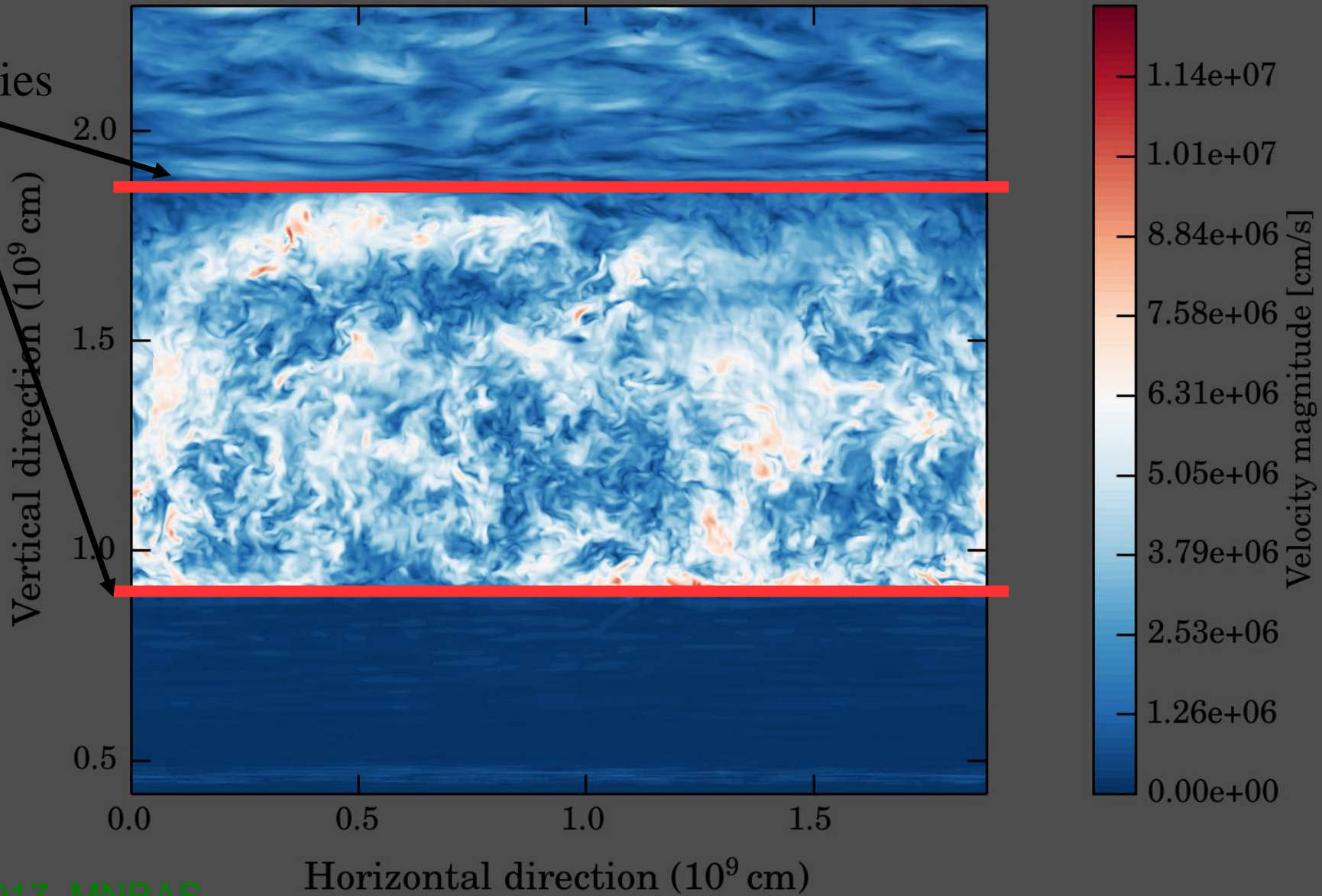


→ Improve theoretical prescriptions

# 3D C-shell Simulations

Snapshot from  $1024^3$  resolution run: Gas Velocity  $\|v\|$

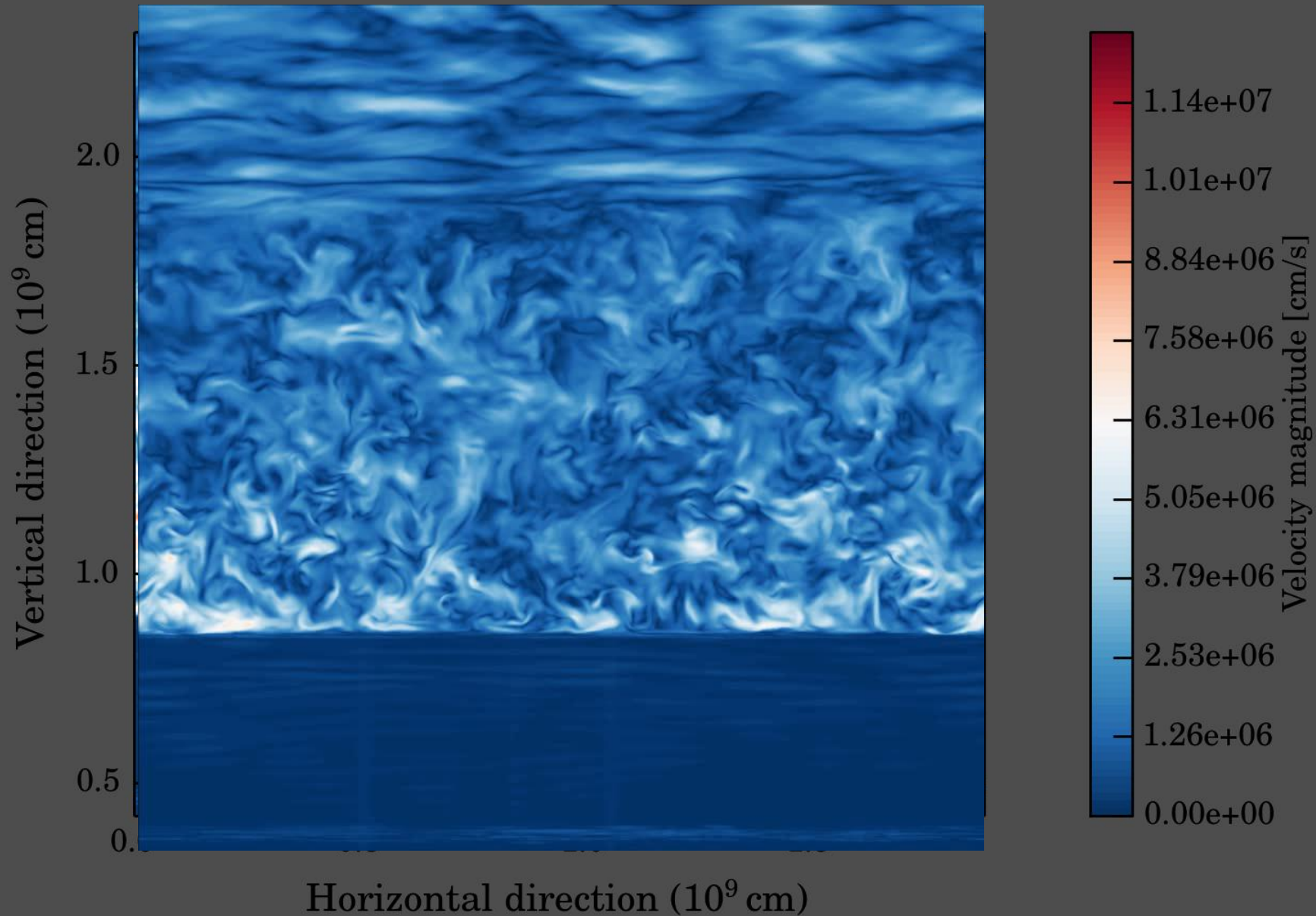
1D boundaries



# 3D C-shell Simulations: $|v|$ movie

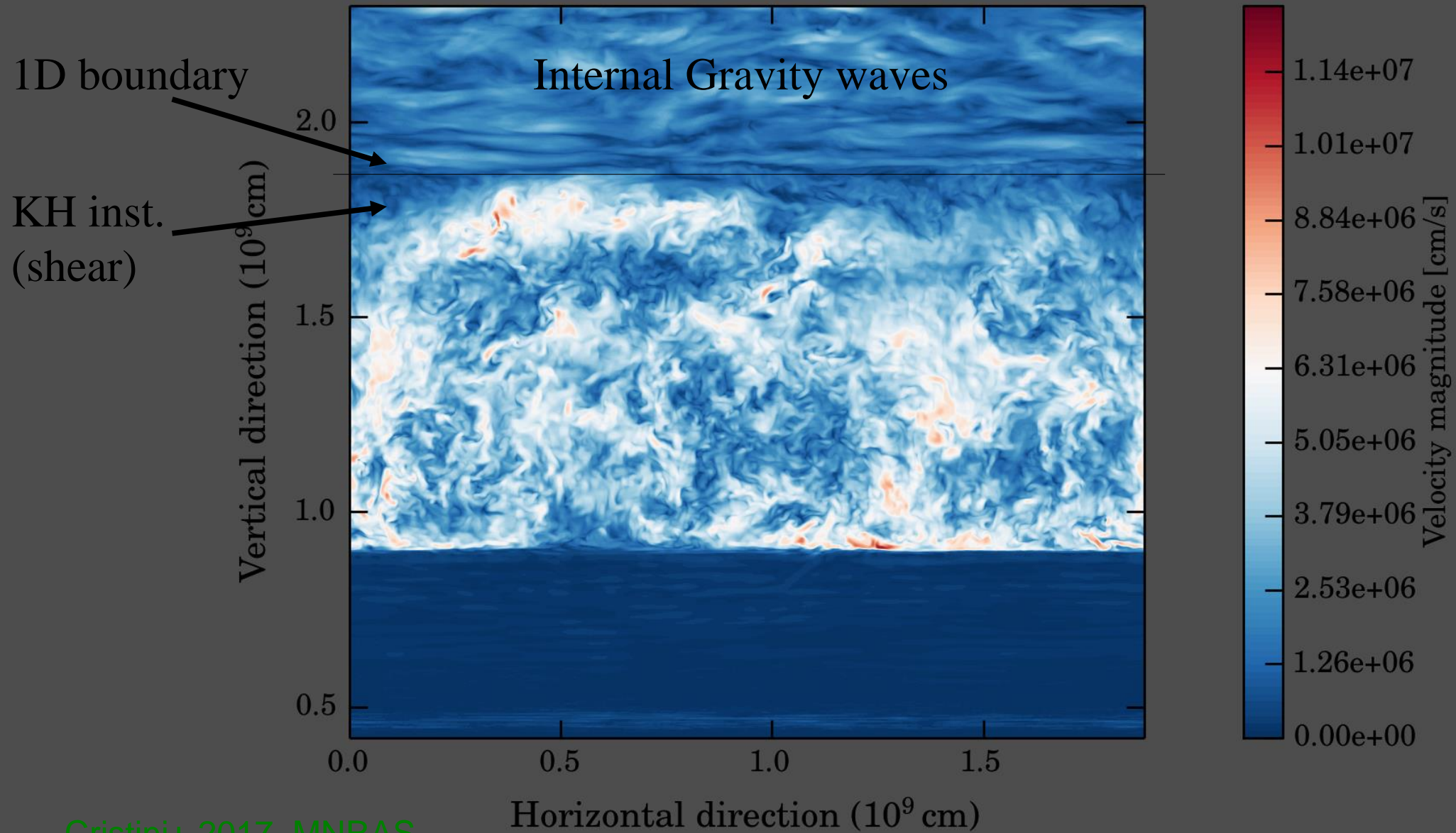
Cristini+ 2017, MNRAS

Gas Velocity  $\|v\|$

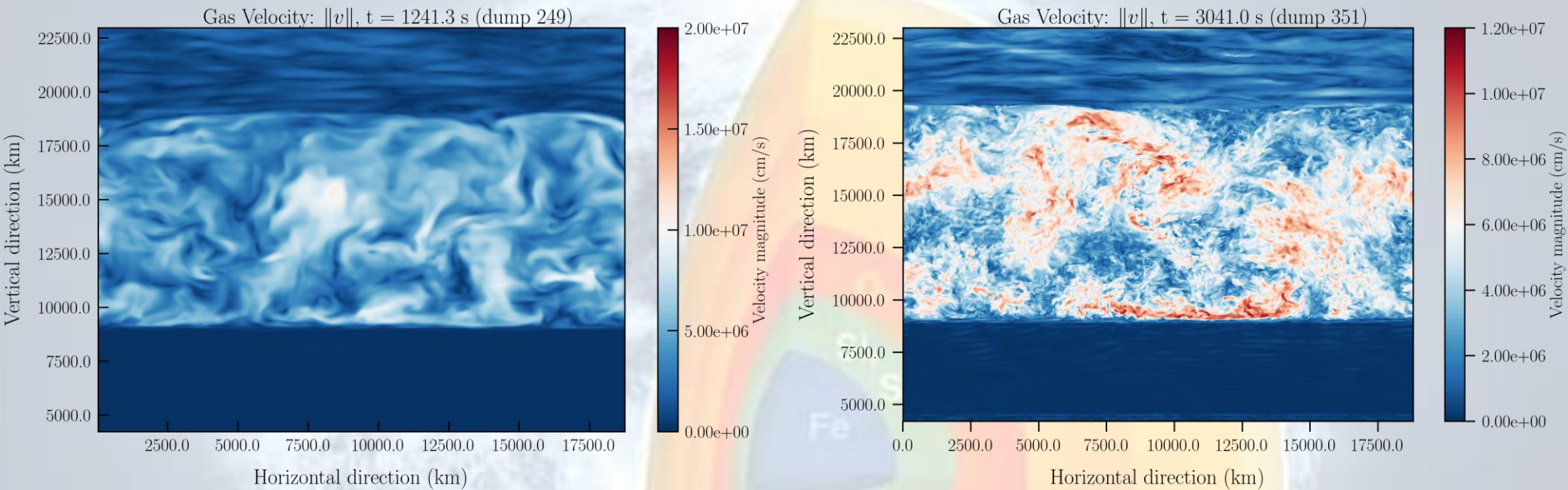


# 3D C-shell Simulations

Snapshot from  $1024^3$  resolution run: Gas Velocity  $\|\mathbf{v}\|$



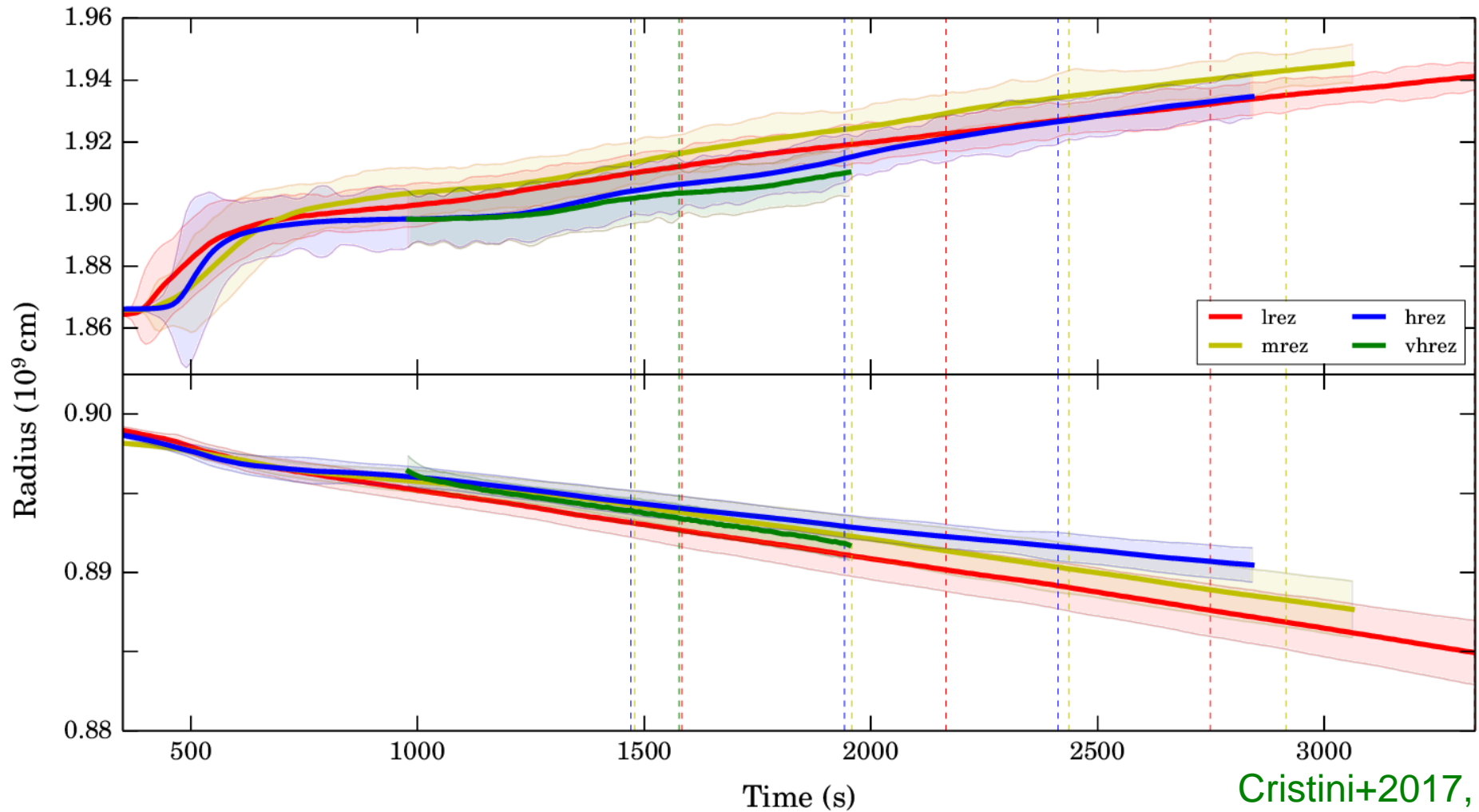
# Resolution & Luminosity Study



Luminosity		128	256	512	1024	
Resolution						
	1			eps1		
	33			eps33		
	100			eps100		
	333			eps333		
	1000	lrez	mrez	hrez/eps1k	vhrez	Resolution study
	3333			eps3k		
	10000			eps10k		
	33333			eps33k		
				Luminosity study		



# Boundary Entrainment



# Entrainment Law Theory

Cristini et al 2019, MNRAS (see also Garcia & Mellado, 2014, Deardor 1980, Chemel, Staquet and Chollet 2010, Fernando, 1991, Stevens and Lenschow, 2001, Jonker+ 2013)

$$\frac{v_e}{v_c} = ARi_B^{-n},$$

A, n are constants

$$Ri_B = \frac{\Delta B \times l}{v_{rms}^2}$$

$$Ri_B = \frac{\text{stabilising potential}}{\text{turbulent kinetic energy}}$$

The buoyancy jump is an integral of the squared buoyancy frequency  $N^2$  with respect to radius  $r$ , given by

$$\Delta b = \int_{r_1}^{r_2} N^2 dr, \quad (2)$$

where  $r_1$  and  $r_2$  encompass the boundary of the convective core, centred at  $r = r_b$ . The upper limit  $r_2$  is equal to  $r_b$  plus some fraction of a pressure scale height; in this study we

A mass entrainment rate,  $\dot{M}_{ent}$ , can be derived from this to give

$$\dot{M}_{ent} = 4\pi r_b^2 \rho_b v_c A R i_B^{-n}, \quad (6)$$

with  $\rho_b$  being the density at  $r = r_b$ . The mass contained within the entrained region,  $M_{ent}$ , is then

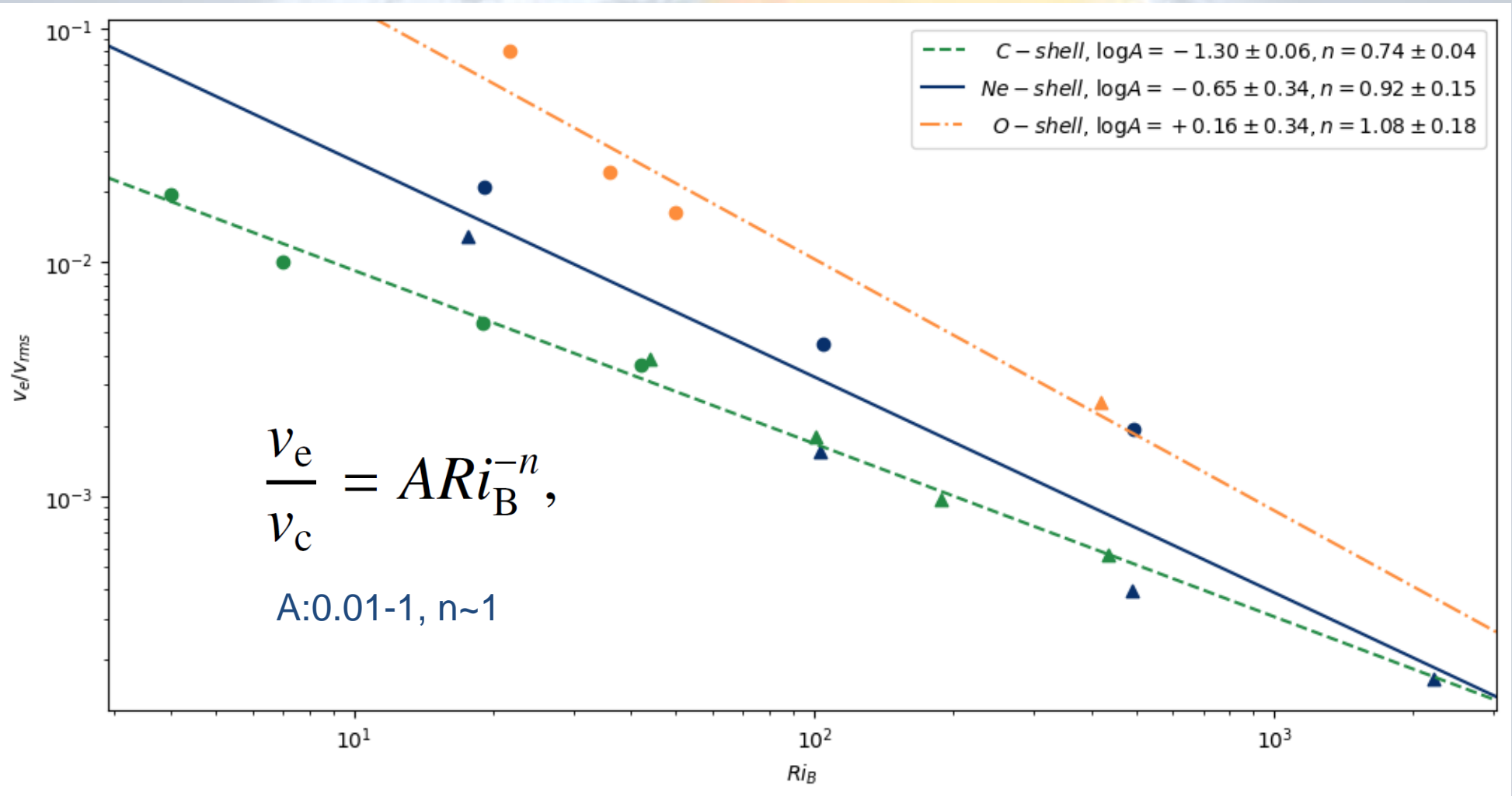
$$M_{ent} = \sum_j \dot{M}_{ent,j} \Delta t_j \quad (7)$$

Instantaneous entrainment as in:  
[Staritsin 2013](#)

Cumulative entrainment as in:  
[Scott+ 2021MNRAS.503.4208S](#)

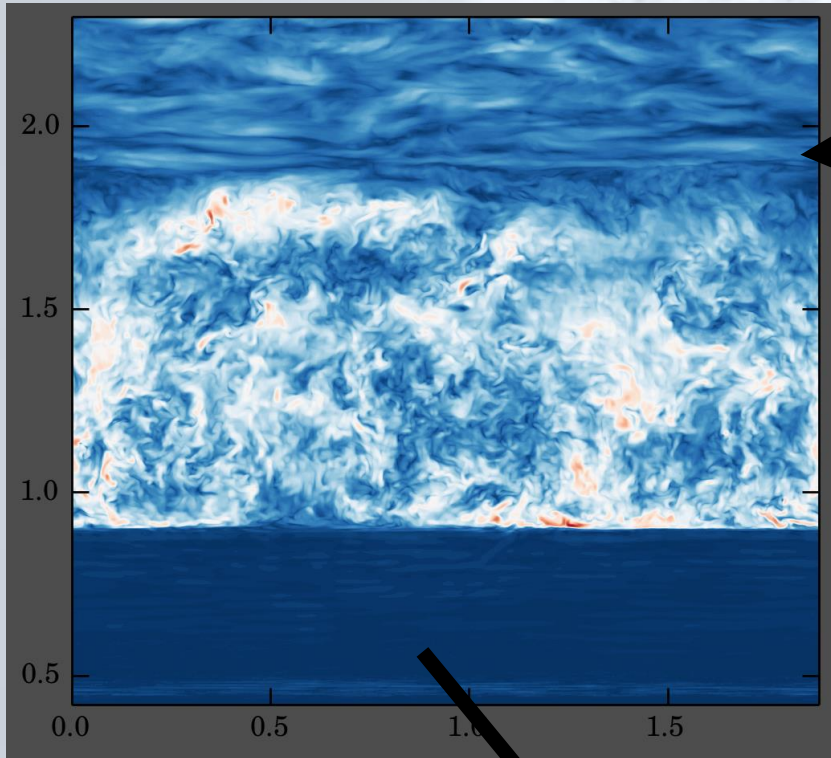
# Entrainment Law for C, Ne & O-burning

Cristini et al 2019, MNRAS (see also Garcia & Mellado, 2014, Deardor 1980, Chemel, Staquet and Chollet 2010, Fernando, 1991, Stevens and Lenschow, 2001, Jonker+ 2013) Georgy+, Rizzuti+, in prep

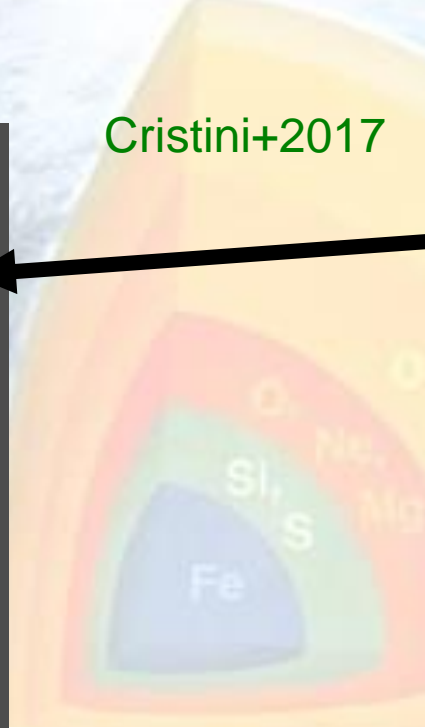


# Way Forward: 1 to 3 to 1D link

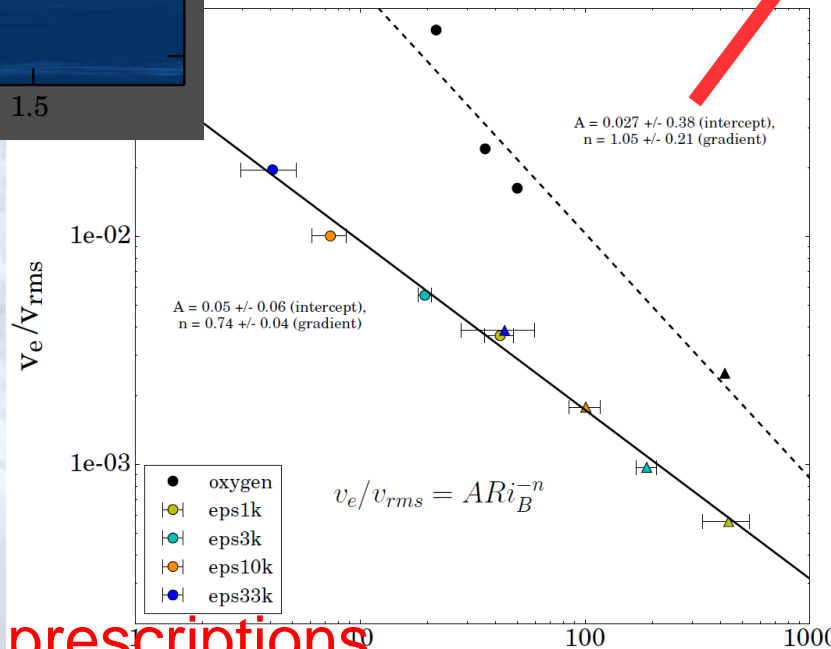
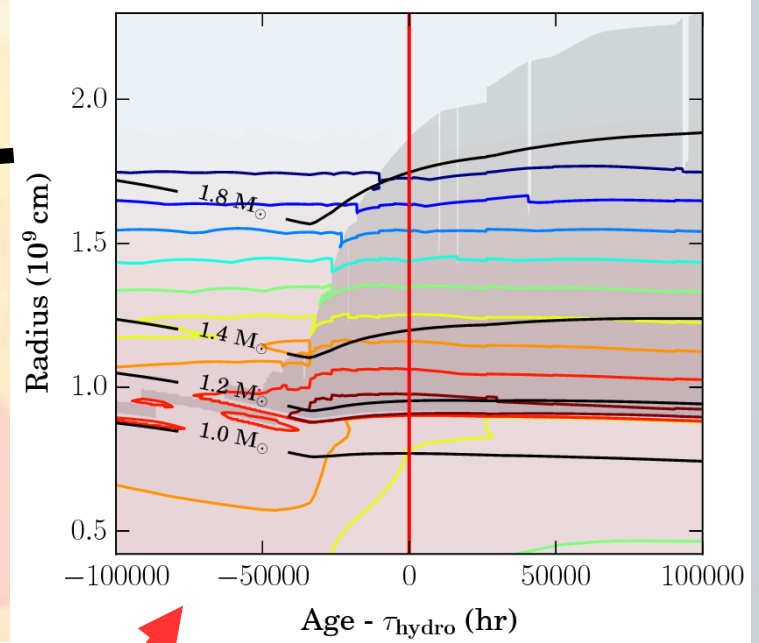
Targeted 3D simulations



Cristini+2017



Uncertainties in 1D



→ Improve theoretical prescriptions

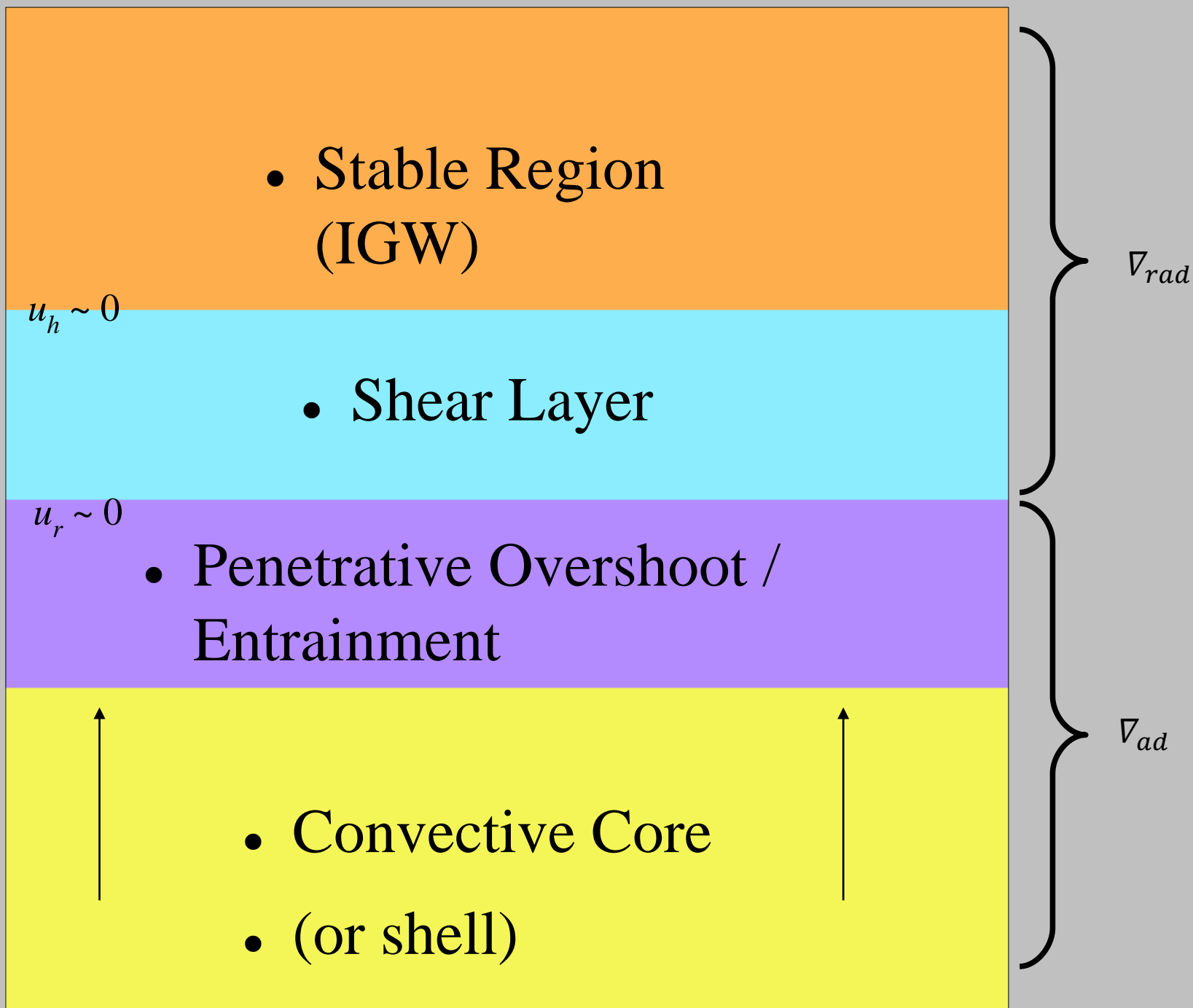
**NEW!** boundary layers



$\nabla_{rad}$

$\nabla_{ad}$

**NEW!** boundary layers

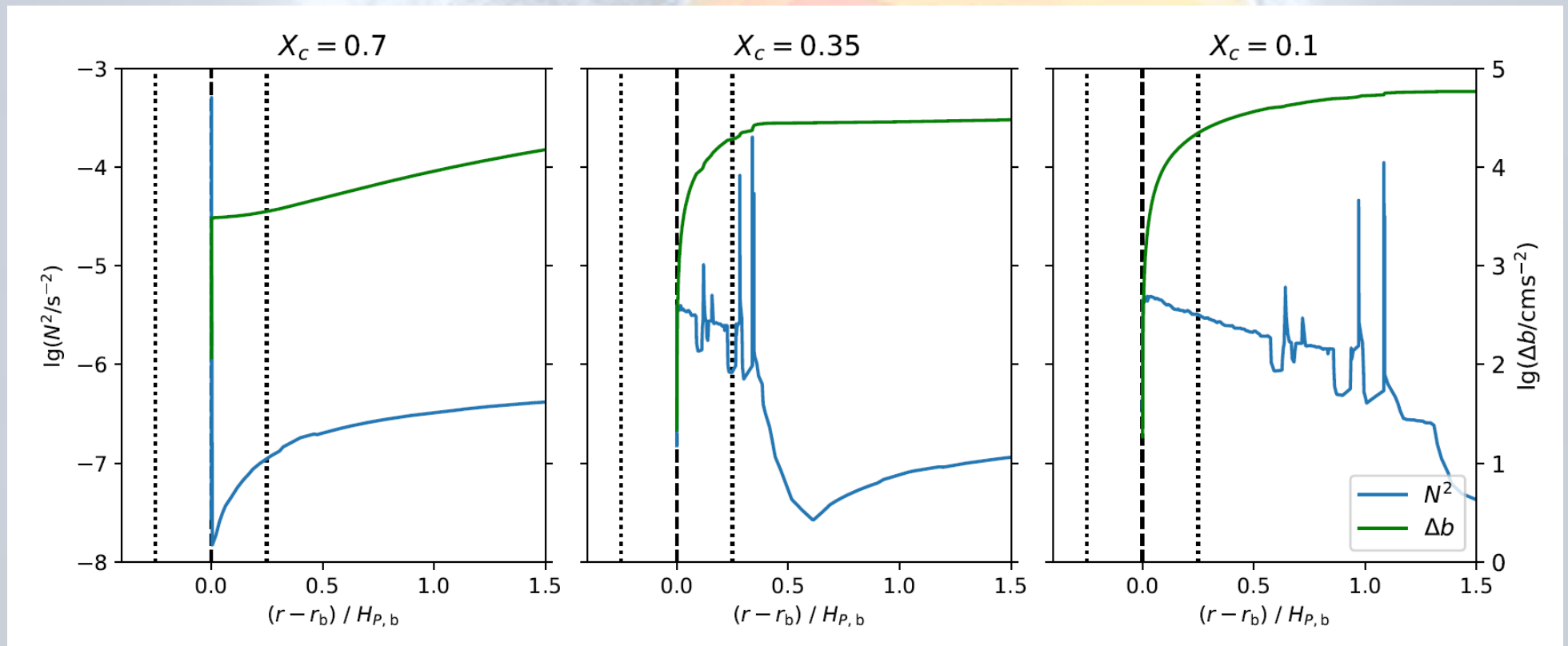


# Implementation of Entrainment in 1D model

Scott+ 2021MNRAS.503.4208S

Calculation of  $\Delta b = \int_{r_1}^{r_2} N^2 dr,$

15  $M_{\odot}$



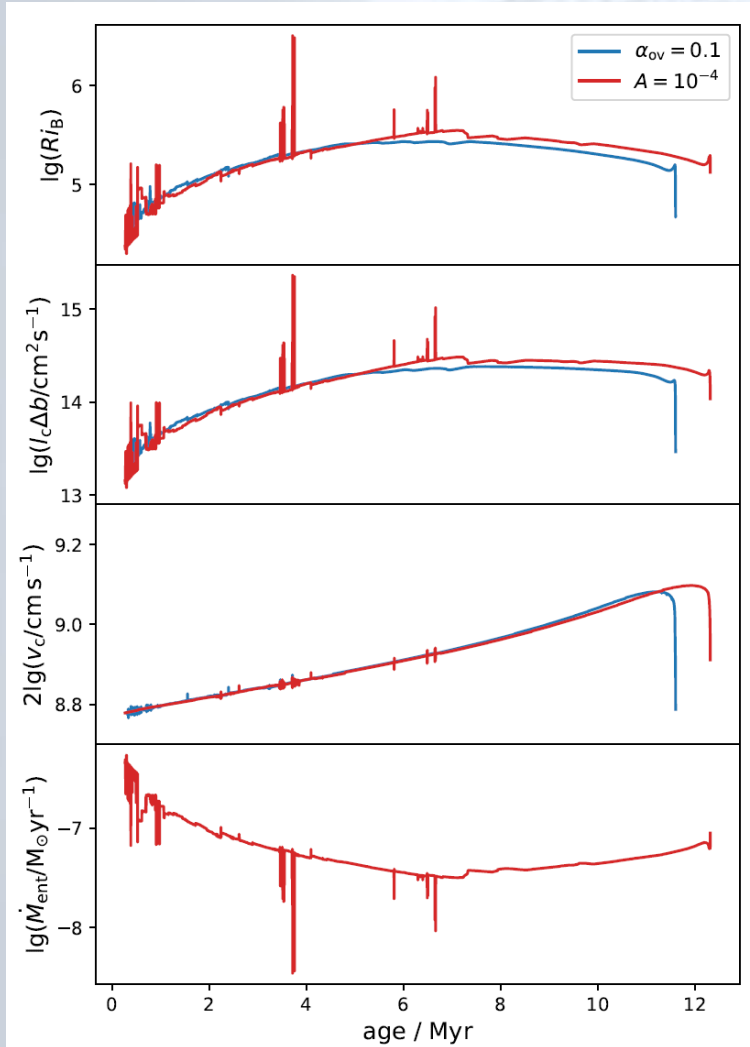
$\Delta b$  increases with time over MS due to  $\mu$ -gradient built-up.

Mild dependence of  $\Delta b$  on choice of  $r_1$  and  $r_2$ , both chosen to be  $H_p/4$

# Implementation of Entrainment in 1D model

Scott+ 2021MNRAS.503.4208S

$15 M_{\odot}$   
 $A=10^{-4}$   
 $n=1$



$$Ri_B = \frac{l_c \Delta b}{v_c^2},$$

$$\dot{M}_{\text{ent}} = 4\pi r_b^2 \rho_b v_c A Ri_B^{-n},$$

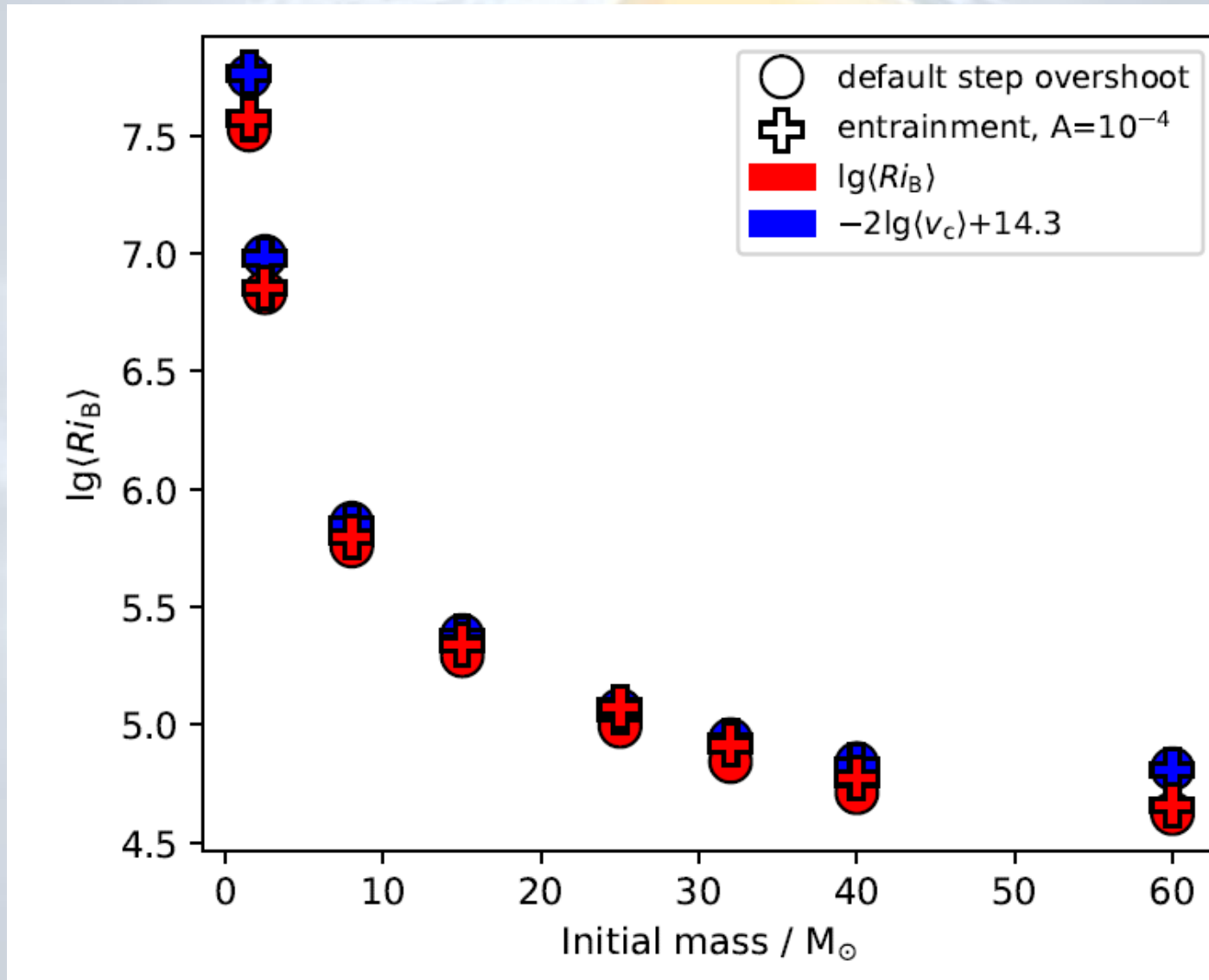
$$M_{\text{ent}} = \sum_j \dot{M}_{\text{ent},j} \Delta t_j$$

RiB increases with time over MS so entrainment rate decreases  
 (small change compared to mass dependence)



# Boundary Penetrability vs Initial Mass

Scott+ 2021MNRAS.503.4208S

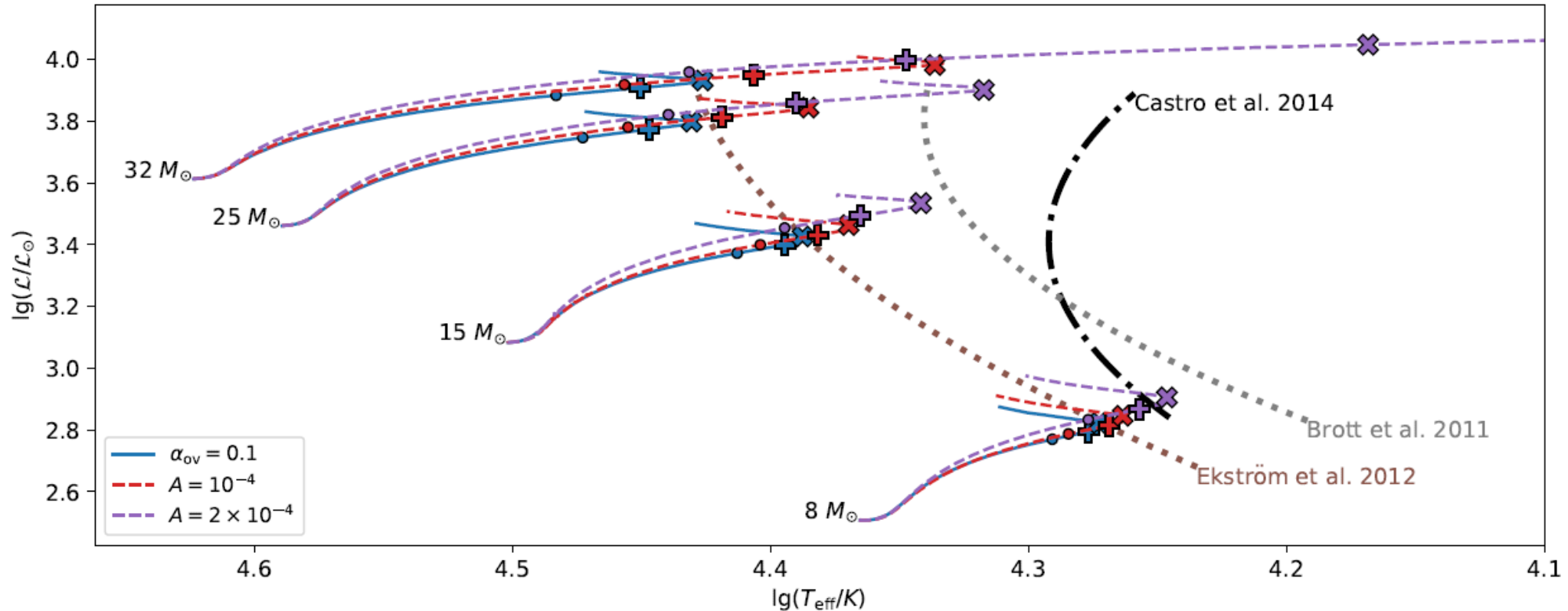


Penetrability ( $\sim 1/Ri_B$ ) increases with initial mass.

Mainly linked to increase in convective velocity,  $v_c$  with mass ( $Ri_B \sim 1/v_c^2$ )

# Entrainment in 1D model: Impact on MS Width

Scott+ 2021MNRAS.503.4208S

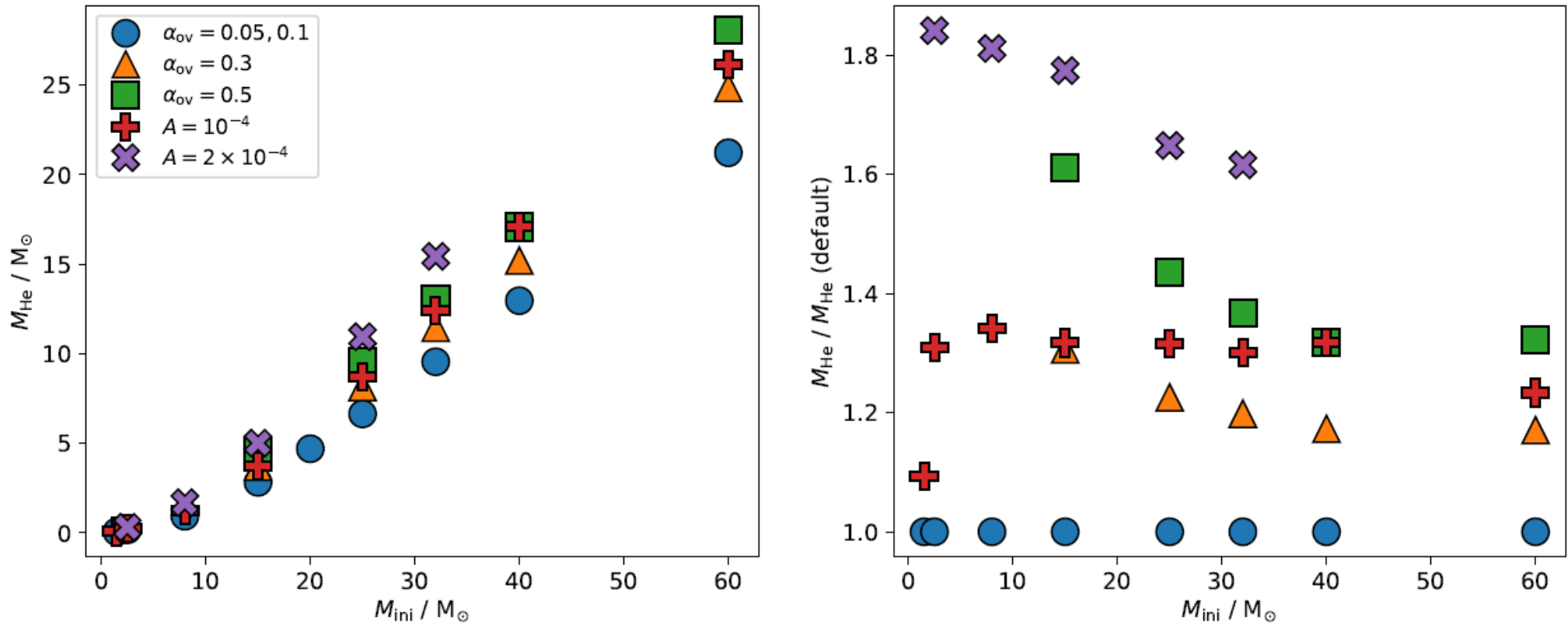


**Figure 6.** Spectroscopic HRD in the mass range 8 to 32  $M_{\odot}$  with  $\alpha_{\text{ov}} = 0.1$  step overshoot and two entrainment models,  $A = 10^{-4}$  and  $2 \times 10^{-4}$ . The dotted lines represent the position of the TAMS from model grids with  $\alpha_{\text{ov}} = 0.1$  (Ekström et al. 2012) and  $\alpha_{\text{ov}} = 0.335$  (Brott et al. 2011). The dash-dotted line represents the position of the empirical TAMS determined by Castro et al. (2014); see their Table 1 for the polynomial coefficients of the three TAMS lines used in this figure. As in Fig. 3, the dots, pluses and crosses have been placed where the model reaches 90%, 95% and 99% of the MS lifetime respectively.

MS width and its mass dependence better reproduced than current models with a single set of parameters

# Entrainment in 1D model: Impact on core masses

Scott+ 2021MNRAS.503.4208S



**Figure 7.** Final helium core mass,  $M_{\text{He}}$ , for various values of step overshoot  $\alpha_{\text{ov}}$  and entrainment parameter  $A$ . All entrainment models use  $n = 1$ . Blue circles represent the default value of  $\alpha_{\text{ov}}$ , which is 0.05 for  $1.5 M_{\odot}$  and 0.1 otherwise. *Left:* Absolute value of  $M_{\text{He}}$  against initial mass. The  $\alpha_{\text{ov}} = 0.1$  point at  $20 M_{\odot}$  is taken from the Ekström et al. (2012) grid. *Right:*  $M_{\text{He}}$  normalised by the default step overshoot value.

CO core masses predicted up to 80% larger than for current models: A  $15 M_{\odot}$  model may behave as a 20-22  $M_{\odot}$  model.

But  $A$  values do not match 3D hydro simulation results yet ...

# Entrainment in 1D model: "Temporary Fix"

Scott+ 2021MNRAS.503.4208S

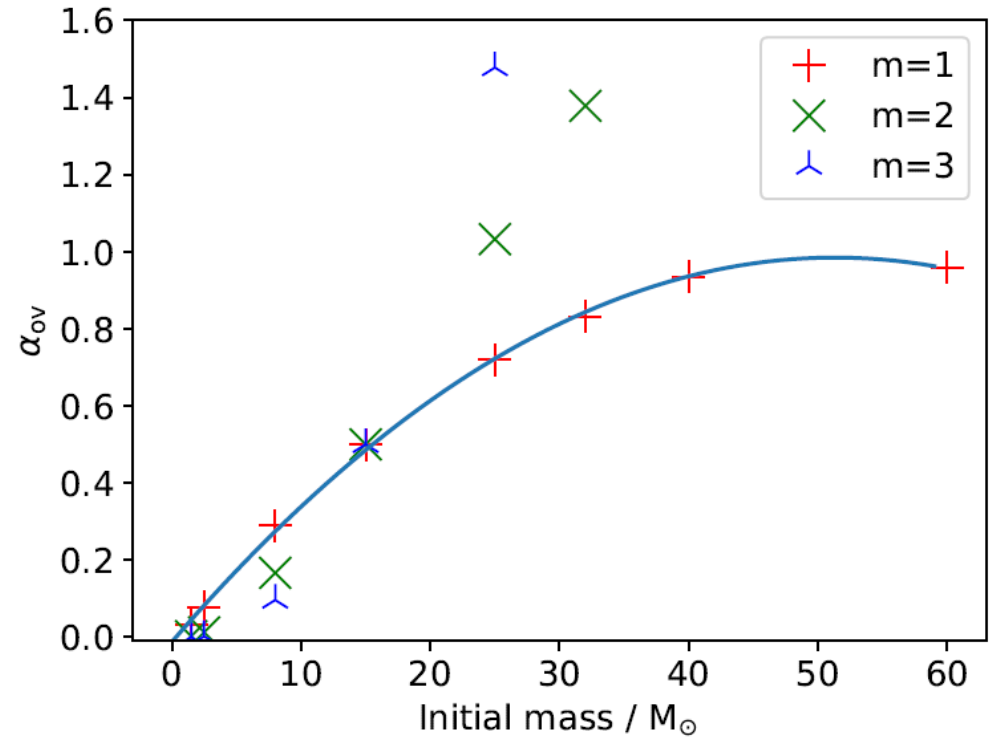
## Entrainment implementation Challenging post-MS:

A mass entrainment rate,  $\dot{M}_{\text{ent}}$ , can be derived from this to give

$$\dot{M}_{\text{ent}} = 4\pi r_b^2 \rho_b v_c A R i_B^{-n}, \quad (6)$$

with  $\rho_b$  being the density at  $r = r_b$ . The mass contained within the entrained region,  $M_{\text{ent}}$ , is then

$$M_{\text{ent}} = \sum_j \dot{M}_{\text{ent},j} \Delta t_j \quad (7)$$

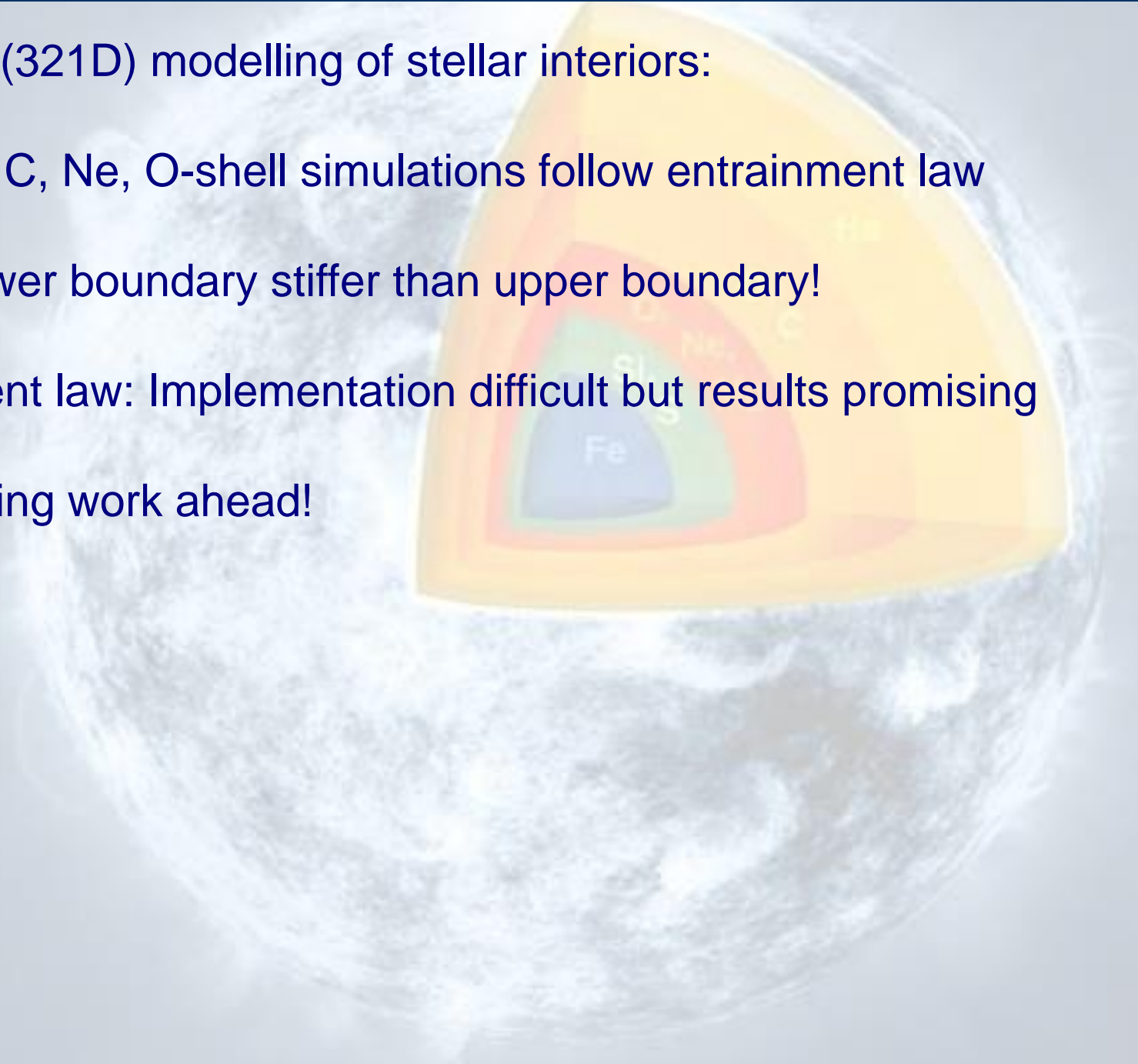


**Figure 9.** Step overshoot parameter  $\alpha_{\text{ov}}$  scaled using Eq.8. The polynomial fit to the  $m = 1$  points uses the equation  $\alpha_{\text{ov}}(M_{\text{ini}}) = -0.00037867M_{\text{ini}}^2 + 0.03885918M_{\text{ini}} - 0.01237484$ . Previous studies such as Claret & Torres (2017) and Moravveji et al. (2016) show that similar results are obtained using an exp-D  $f$  parameter which is roughly a factor of 10 to 15 smaller than the equivalent step overshoot  $\alpha_{\text{ov}}$ . Therefore a fit for  $f(M_{\text{ini}})$  would be roughly 1/10 to 1/15 of  $\alpha_{\text{ov}}(M_{\text{ini}})$ .

Mass-dependent  $\alpha_{\text{ov}}/f_{\text{ov}}$  (fits given in caption)

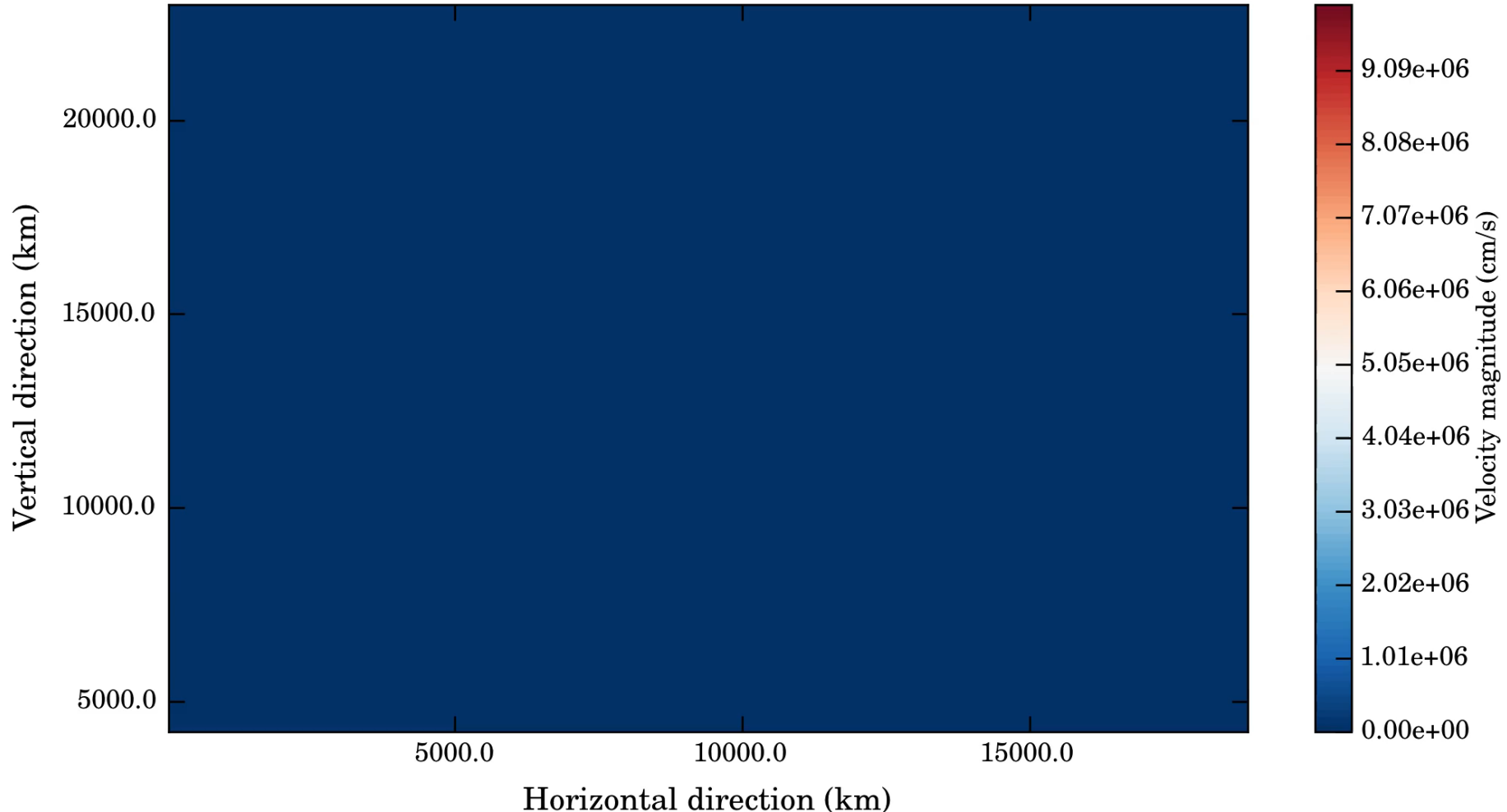
# Conclusions & Outlook

- 1 - 3 - 1 D (321D) modelling of stellar interiors:
  - 3D C, Ne, O-shell simulations follow entrainment law
  - Lower boundary stiffer than upper boundary!
- Entrainment law: Implementation difficult but results promising
- More exciting work ahead!



# 2D Simulations of Convection

Gas Velocity:  $\|v\|$ ,  $t = 0.0$  s (dump 1)



2D simulations a factor of  $N_z$  cheaper so 100 – 1000 times cheaper but ...

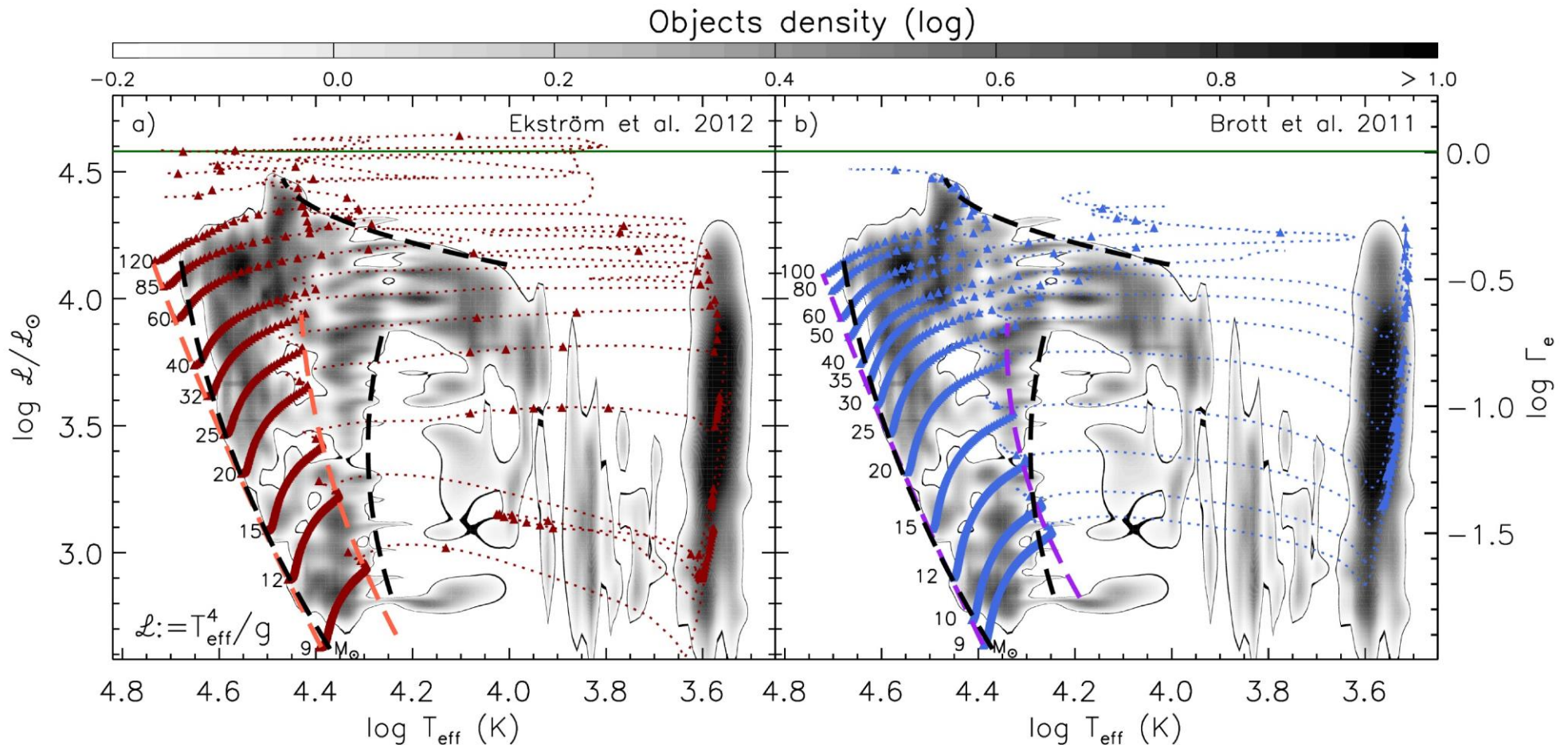
TKE cascade goes upwards rather than downwards!

Vortices survive much longer and merge rather than decay!

Convection is 3D and turbulent so 2D simulations may never reproduce real flow!!

# 1D Model Uncertainties: MS Width

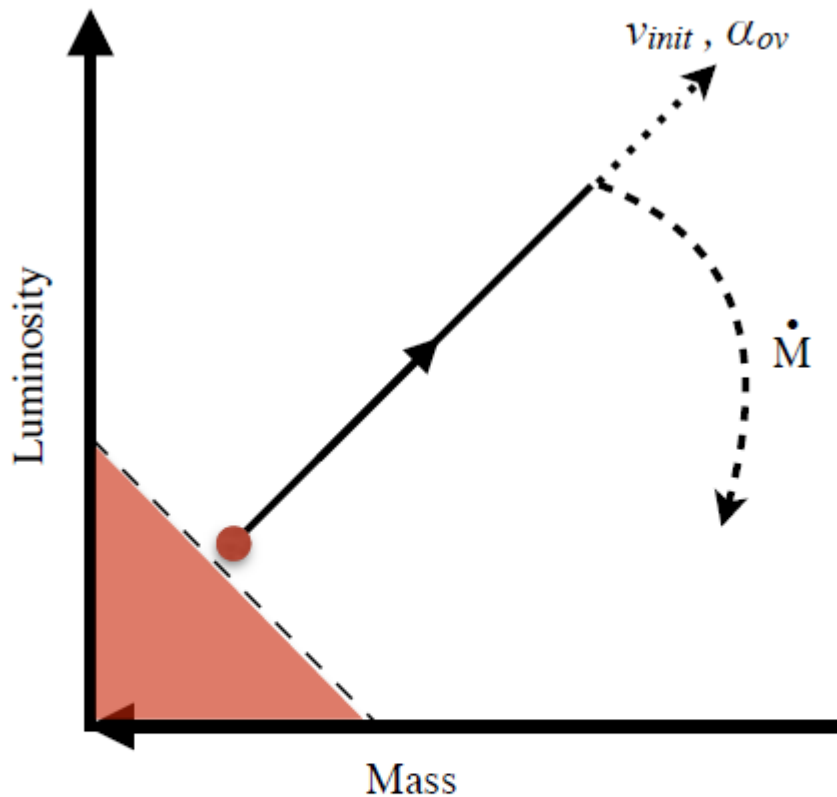
sHRD: Castro et al (2014A&A...570L..13C),  
Langer & Kudritzki (2014, A&A, 564, A52)



MS width and its mass dependence not reproduced by current models

(but see Vink et al. 2010; McEvoy et al. 2015)

# 1D Model Uncertainties: L-M Plane



**Fig. 4.** Illustration of the mass-luminosity plane with a typical evolutionary track entering the ZAMS at the red dot, evolving along the black arrow. The dotted vector suggests how increased rotation and/or convective overshooting may extend the  $M-L$  vector. The curved dashed line represents the gradient at which mass-loss rates affect this  $M-L$  vector. The red solid region represents the boundary set by the mass-luminosity relationship, and as such is forbidden.

Higgins & Vink  
(A&A 622, A50, 2019)

$\alpha_{ov}$  up to  $0.5 H_p$  and  
rotation-induced mixing  
both needed to explain  
HD166734

See also Martinet et al 2021, A&A

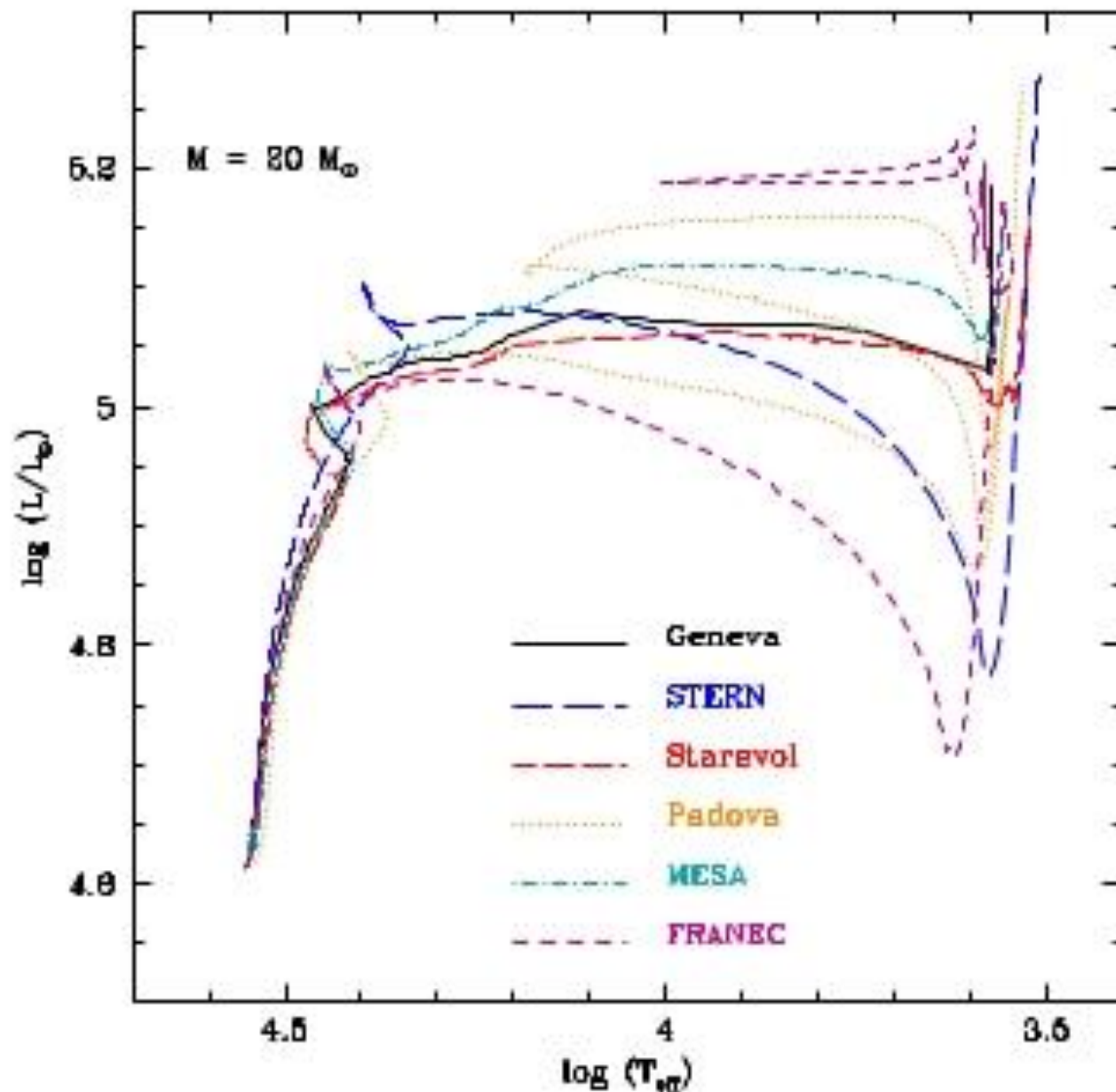
$\alpha_{ov}$  up to  $0.5 H_p$  ;  $f_{ov}$  up to 0.05 from Asteroseismology

See e.g. Deheuvels et al. 2016, Aerts et al 2018, Claret & Torres 2019, Pedersen et al 2020, 2021; Aerts et al ...



# 1D Model Uncertainties: Post-MS

*Martins and Palacios (2013)*



Different prescriptions for convective mixing and free parameters **strongly affect** post-MS evolution.

See also [Georgy+2014, A&A](#)

[Jones et al 2015, MNRAS, 447, 3115](#)

# 1D Model Uncertainties: Core Masses

Kaiser et al, MNRAS 496, 1967 (2020)

**Table 2.** The absolute and relative variation of the total mass,  $M_{\text{tot}}$ , the helium core mass,  $M_{\alpha}$ , and the carbon-oxygen core mass,  $M_{\text{CO}}$ . The individual values of each model are shown in Table 1. The values include *Ledoux* and *Schwarzschild* models.

	$\Delta M_{\text{tot}} (M_{\odot})$	$\delta M_{\text{tot}}$	$\Delta M_{\alpha} (M_{\odot})$	$\delta M_{\alpha}$	$\Delta M_{\text{CO}} (M_{\odot})$	$\delta M_{\text{CO}}$
15 $M_{\odot}$	3.38 (3.45)	27.72% (27.80%)	2.30 (2.74)	43.07% (51.31%)	2.40 (3.94)	69.77% (114.53%)
20 $M_{\odot}$	8.60	74.35%	2.94	36.70%	2.94	51.49%
25 $M_{\odot}$	9.02	65.84%	3.96	40.08%	4.1	55.48%

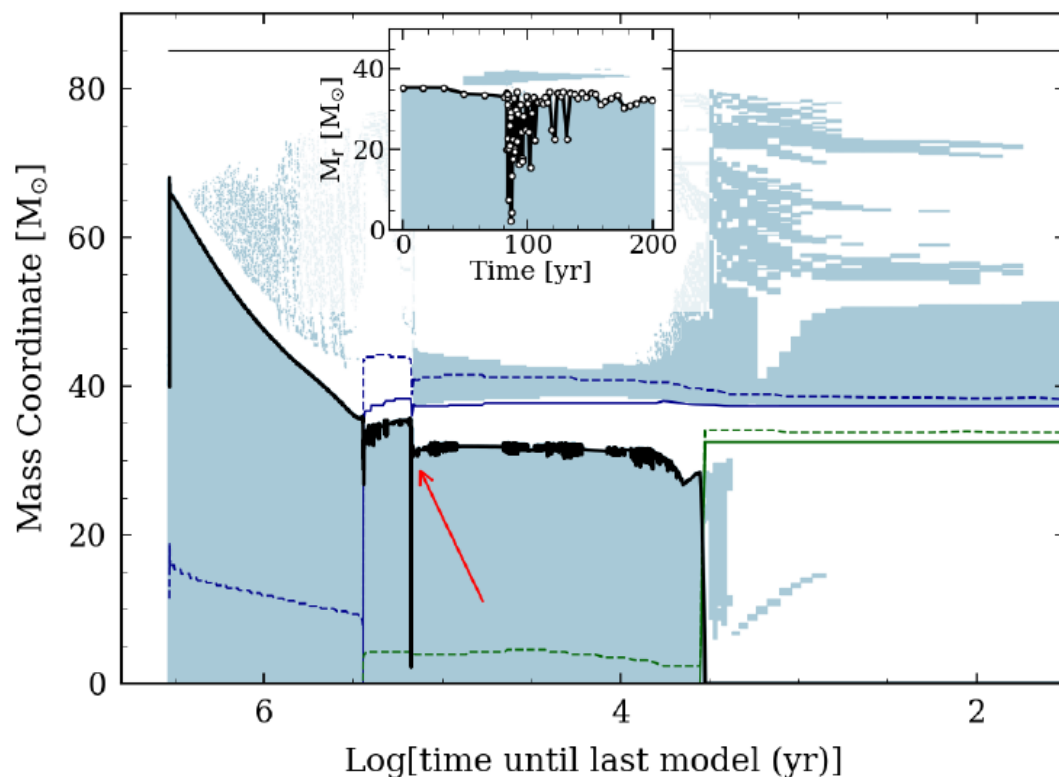
*Notes.* The variation for a quantity  $Q$  is calculated as  $\delta Q = \frac{\Delta Q}{Q_{\text{ref}}} \equiv \frac{Q_{\text{max}} - Q_{\text{min}}}{Q_{\text{ref}}} \times 100$ , where  $Q_{\text{max}}$  and  $Q_{\text{min}}$  are the maximal and minimal values of the quantity for the initial mass and  $Q_{\text{ref}}$  the value of the reference model (see the text).

\* The values in parentheses include the models with no CBM.

Up to 70% uncertainties in CO core masses predicted by models using the full range of CBM parameters. A 15  $M_{\odot}$  model may behave as a 20  $M_{\odot}$  model.

# 1D Model Uncertainties: BH-mass gap

*e.g. Farrel et al, MNRAS (2020) & references therein*



Shell mergers/interactions  
reduce CO-core mass at  
very low Z.

**Figure 2.** Kippenhahn diagram of a GENEC non-rotating  $85 M_{\odot}$  model at  $Z = 0$ . Solid (dashed) lines correspond to the peak ( $100 \text{ erg/g/s}$ ) of the energy generation rate for H burning (blue) and He burning (green). The red arrow indicates the H-He shell interaction. An inset is included at the top of the figure to show that the interaction is resolved, where white circles indicate each timestep.

# 1D Model Uncertainties: Post-MS

Detailed convective shell history affects fate of models: strong/weak/failed explosions!!!

Sukhbold & Woosley, 2014ApJ...783...10S

Sukhbold, Ertl et al, 2016ApJ...821...38S,

Ugliano et al 2012, Ertl et al 2015

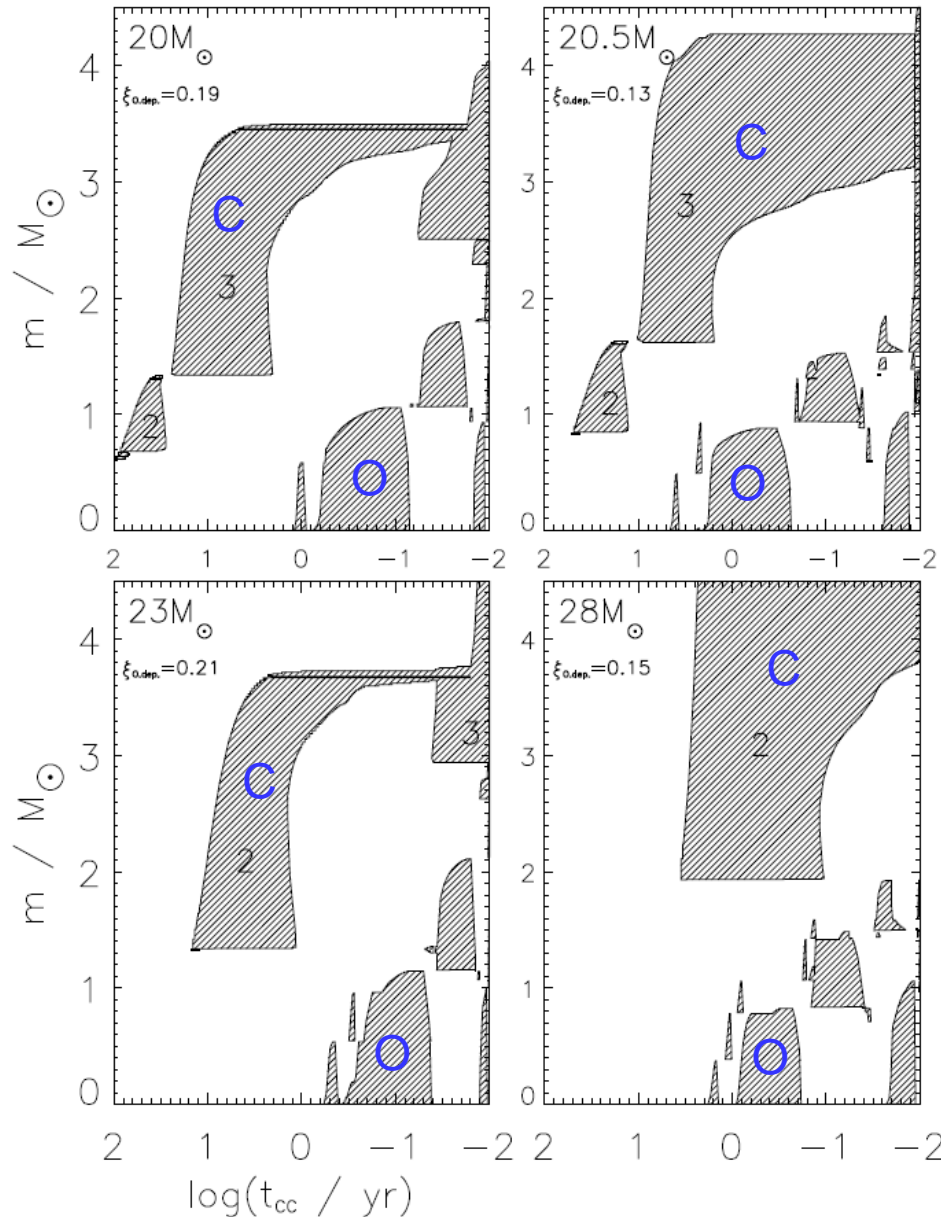


FIG. 13.— Convective history of four models showing the major

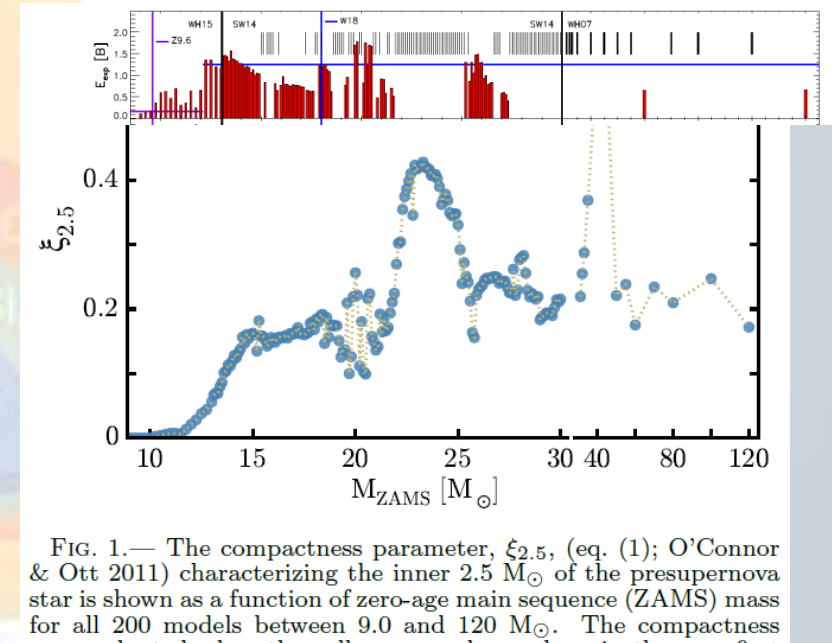


FIG. 1.— The compactness parameter,  $\xi_{2.5}$ , (eq. (1); O'Connor & Ott 2011) characterizing the inner  $2.5 M_{\odot}$  of the presupernova star is shown as a function of zero-age main sequence (ZAMS) mass for all 200 models between  $9.0$  and  $120 M_{\odot}$ . The compactness

## Non-monotonic behaviour?

We are particularly interested in how the “explodability” of the presupernova models and their observable properties correlate with their “compactness” (Fig. 1; O'Connor & Ott 2011)

$$\xi_M = \frac{M/M_{\odot}}{R(M)/1000 \text{ km}} \Big|_{t_{\text{bounce}}}, \quad (1)$$

and other measures of presupernova core structure (§ 3.1.3; Ertl et al. (2015)). Using a standard central engine in presupernova models of variable compactness, a significant correlation in outcome is found (§ 4). As pre-